

215(1): Interpretation of Planetary Precession as a Coriolis Force

Consider the fundamental theorem:

$$\left(\frac{d\underline{Q}}{dt}\right)_{\text{fixed}} = \left(\frac{d\underline{Q}}{dt}\right)_{\text{rotating}} + \underline{\omega} \times \underline{Q} \quad - (1)$$

where \underline{Q} is any vector. Here \underline{Q} subscript refers to fixed and rotating frame. Let:

$$\underline{Q} = \underline{p} \quad - (2)$$

$$\text{where } \underline{p} = m\underline{v} \quad - (3)$$

is total linear momentum. Then:

$$\underline{F}_{\text{fixed}} = \underline{F}_{\text{rotating}} + m \underline{\omega} \times \underline{v} \quad - (4)$$

$$\text{where } \underline{F}_{\text{Coriolis}} = m \underline{\omega} \times \underline{v} \quad - (5)$$

is the Coriolis force.

As in UFT 214:

$$\underline{F}_{\text{fixed}} = -\frac{k}{r^2} \left(x^2 + (1-x^2) \right) \frac{d}{r} \underline{e} \quad - (6)$$

$$\underline{F}_{\text{rotating}} = -\frac{k}{r^2} \underline{e} \quad - (7)$$

where

2)

$$\underline{e}_{rf} = \underline{i} \cos \theta + \underline{j} \sin \theta \quad - (8)$$

$$\underline{e}_{rot} = \underline{i} \cos(x\theta) + \underline{j} \sin(x\theta) \quad - (9)$$

$$k = \frac{mMG}{L^2} \quad - (10)$$

$$\alpha = \frac{L^2}{mk} \quad - (11)$$

So :

$$\underline{F}_{Coriolis} = \underline{F}_{fixed} - \underline{F}_{rotating} \quad - (12)$$

$$= -\frac{k}{r^2} \left(x^2 + (1-x^2) \frac{\alpha}{r} \right) \left(\underline{i} \cos \theta + \underline{j} \sin \theta \right)$$

$$+ \frac{k}{r^2} \left(\underline{i} \cos(x\theta) + \underline{j} \sin(x\theta) \right)$$

$$= -\frac{k}{r^2} \left[\underline{i} \left(\left(x^2 + (1-x^2) \frac{\alpha}{r} \right) \cos \theta - \cos(x\theta) \right) + \underline{j} \left(\left(x^2 + (1-x^2) \frac{\alpha}{r} \right) \sin \theta - \sin(x\theta) \right) \right] \quad - (13)$$

$$x = 1 \quad - (14)$$

When

der

$$\underline{F}_{Coriolis} = \underline{0} \quad - (15)$$