

## 259(1): Electrostatics in ECE Theory

The equations of electrostatics in ECE theory are the Coulomb law:

$$\underline{\nabla} \cdot \underline{E}^a = \underline{\omega}^a_b \cdot \underline{E}^b - c \underline{A}^b \cdot \underline{R}^a_b(ad), - (1)$$

The Faraday law of induction in the absence of a magnetic field:

$$\underline{\nabla} \times \underline{E}^a = \underline{0} - (2)$$

and definition of the electric field:

$$\underline{E}^a = -c \underline{\nabla} A^a_0 - \frac{\partial \underline{A}^a}{\partial t} - c \underline{\omega}^a_b A^b + c A^b_0 \underline{\omega}^a_b - (3)$$

In electrostatics the Ampère Maxwell law reduces to:

$$\frac{\partial \underline{E}^a}{\partial t} = \underline{0} - (4)$$

In the absence of a magnetic field:

$$\underline{B}^a = 0, \underline{A}^a = 0, \underline{J}^a = \underline{0} - (5)$$

are possible solutions. However the Aharonov Bohm effects we described in the absence of a magnetic flux density

Sy: 
$$\underline{\nabla} \times \underline{A}^a = \underline{\omega}^a_b \times \underline{A}^b - (6)$$

$$\underline{\nabla} \times \underline{A}^a = \kappa \underline{A}^a - (7)$$

where:

and:

$$\underline{\nabla} \times \underline{\omega}^a_b = -i\kappa \underline{\omega}^a_b \quad - (8)$$

The Aharonov Bohm effects are therefore described by:

$$\kappa \underline{A}^a = \underline{\omega}^a_b \times \underline{A}^b \quad - (9)$$

The electric charge density and electric current density are described in ECE theory by:

$$\rho^a = \epsilon_0 (\underline{\omega}^a_b \cdot \underline{E}^b - c \underline{A}^b \cdot \underline{R}^a_b(\text{orb})) \quad - (10)$$

and:

$$\underline{J}^a = \epsilon_0 c (\underline{\omega}^a_b \cdot \underline{E}^b - c \underline{A}^b \cdot \underline{R}^a_b(\text{orb}) + c \underline{\omega}^a_b \times \underline{B}^b - c \underline{A}^b \times \underline{R}^a_b(\text{spin})) \quad - (11)$$

The correct S.I. units are:

$$E = \text{Vm}^{-1} = \text{JC}^{-1}\text{m}^{-1}$$

$$A = \text{JsC}^{-1}\text{m}^{-1}$$

$$\epsilon_0 = \text{J}^{-1}\text{C}^2\text{m}^{-1}$$

$$B = \text{JsC}^{-1}\text{m}^{-2} = \text{tesla}$$

$$\rho = \text{Cm}^{-3}$$

$$\underline{J} = \text{Cm}^{-2}\text{s}^{-1}$$

$$\omega = \text{m}^{-1}$$

$$R = \text{m}^{-2}$$

So it may be checked that the above equations

→ (2) and (3) have the right units.

In the absence of a vector potential:

$$\underline{\nabla} \cdot \underline{E}^a = \underline{\omega}^a_b \cdot \underline{E}^b - (12)$$

and

$$\underline{J}^a = \underline{0} = \epsilon_0 c \left( \underline{\omega}^a_b \underline{E}^b - c A^b \cdot \underline{R}^a_b(\omega b) \right) - (13)$$

so

$$\underline{\omega}^a_b \underline{E}^b = c A^b \cdot \underline{R}^a_b(\omega b) - (14)$$

and

$$\underline{E}^a = -c \underline{\nabla} A^a_0 + c A^b \cdot \underline{\omega}^a_b - (15)$$

w/:

$$\underline{\nabla} \times \underline{E}^a = \underline{0} - (16)$$

From eqs. (15) and (16):

$$\underline{\nabla} \times \underline{E}^a = -c \underline{\nabla} \times \underline{\nabla} A^a_0 + c \underline{\nabla} \times (A^b \cdot \underline{\omega}^a_b) - (17)$$

so

$$\boxed{\underline{\nabla} \times (A^b \cdot \underline{\omega}^a_b) = 0} - (18)$$

in electrostatics. This is because:

$$\underline{\nabla} \times \underline{\nabla} A^a_0 = 0 - (19)$$

4) Eq. (18) is :

$$\underline{\nabla} \times (\underline{A}^b \cdot \underline{\omega}^a_b) = \underline{A}^b \cdot \underline{\nabla} \times \underline{\omega}^a_b + (\underline{\nabla} \underline{A}^b) \times \underline{\omega}^a_b = 0 \quad - (20)$$

where  $\phi^a = c A^a_0 \quad - (21)$

i.e. scalar potential. So:

$$\phi^a \underline{\nabla} \times \underline{\omega}^a_b + (\underline{\nabla} \phi^b) \times \underline{\omega}^a_b = 0 \quad - (22)$$

The magnetic charge density in ERI theory is given by:

$$\rho_{\text{mag}}^a = \epsilon_0 c (\underline{\omega}^a_b \cdot \underline{B}^b - \underline{A}^b \cdot \underline{R}^a_b(\text{spin})) \quad - (23)$$

and the magnetic current density is:

$$\underline{J}_{\text{mag}}^a = \epsilon_0 (\underline{\omega}^a_b \times \underline{E}^b - c \omega^a_{0b} \underline{B}^a - c (\underline{A}^b \times \underline{R}^a_b(\text{orb}) - \underline{A}^b \cdot \underline{R}^a_b(\text{spin}))) \quad - (24)$$

These are thought to vanish in the theory of electromagnetism. So:

o) A rich variety of phenomena come out of these equations.

The Coulomb Law

This may be expressed along the  $z$  axis as:

$$\underline{E}^{(3)} = - \frac{e}{4\pi\epsilon_0 z^2} \underline{e}^{(3)} \quad - (37)$$

where

$$\underline{E} = \underline{E}^{(3)} = E_z \underline{k}, \quad - (38)$$

$$\underline{e}^{(3)} = \underline{k} \quad - (39)$$

So: 
$$\underline{\nabla} \cdot \underline{E} = \frac{\partial E_z}{\partial z} = \frac{e}{2\pi\epsilon_0 z^3} \quad - (40)$$

There is only one sense of polarization, so:

$$(a) = (b) = (3) \quad - (41)$$

and 
$$\frac{\partial E_z}{\partial z} = \omega E_z = \omega z^{(3)} \cdot \underline{E}^{(3)} \quad - (42)$$

$$\omega := \omega z^{(3)} \quad - (43)$$

where

$$E_z = - \frac{e}{4\pi\epsilon_0 z^2} \quad - (44)$$

therefore

$$\frac{e}{2\pi\epsilon_0 z^3} = - \omega \frac{e}{4\pi\epsilon_0 z^2} \quad - (45)$$

and

$$\boxed{\omega = - \frac{2}{z}} \quad - (46)$$

$$\underline{\omega}^a{}_b \cdot \underline{B}^b = \underline{A}^b \cdot \underline{R}^a{}_b(\text{spin}) - (25)$$

and:

$$\begin{aligned} \underline{\omega}^a{}_b \times \underline{E}^b - c \underline{\omega}^a{}_b \underline{B}^a - c \underline{A}^b \times \underline{R}^a{}_b(\text{orb}) + \underline{A}^b \cdot \underline{R}^a{}_b(\text{spin}) \\ = \underline{0} \end{aligned} - (26)$$

in which:

$$\underline{\nabla} \times \underline{B}^a = \kappa \underline{B}^a - (27)$$

$$\underline{\nabla} \times \underline{A}^a = \kappa \underline{A}^a - (28)$$

In electrostatics, eq. (25) is true because:

$$\underline{B}^b = \underline{0} - (29)$$

$$\underline{A}^b = \underline{0} - (30)$$

In electrostatics eq. (26) becomes: - (31)

$$\underline{\omega}^a{}_b \times \underline{E}^b + \underline{A}^b \cdot \underline{R}^a{}_b(\text{spin}) = \underline{0}$$

Summary of Equation of Electrostatics

$$\underline{\nabla} \cdot \underline{E}^a = \underline{\omega}^a{}_b \cdot \underline{E}^b - (32)$$

$$\underline{\omega}^a{}_b \underline{E}^b = \phi^b \underline{R}^a{}_b(\text{orb}) - (33)$$

$$c \underline{\omega}^a{}_b \times \underline{E}^b + \phi^b \underline{R}^a{}_b(\text{spin}) = \underline{0} - (34)$$

$$\phi^a \underline{\nabla} \times \underline{\omega}^a{}_b + (\underline{\nabla} \phi^b) \times \underline{\omega}^a{}_b = \underline{0} - (35)$$

$$\underline{E}^a = -\underline{\nabla} \phi^a + \phi^b \underline{\omega}^a{}_b - (36)$$

In general :

$$\underline{\nabla} \times \underline{\omega}^a b = 1 \tau \underline{\omega}^a b - (47)$$

as derived in 4FT 258, but if it is assumed that the only relevant spin connection is :

$$\underline{\omega} = \omega \underline{k} = -\frac{2}{Z} \underline{k} - (48)$$

then :

$$\boxed{\underline{\nabla} \times \underline{\omega} = \underline{0}} - (49)$$

This is a Beltrami equation of the type :

$$\boxed{\underline{\nabla} \times \underline{\omega} = 0 \underline{\omega}} - (50)$$

W. of the assumption (48) it follows that :

$$\underline{\omega} \times \underline{E} = \underline{0} - (51)$$

so

$$\underline{R}(\text{spin}) = \underline{0} - (52)$$

from eq. (33). Self consistently, from eq. (35):

$$(\underline{\nabla} \phi) \times \underline{\omega} = \underline{0} - (53)$$

From eqs. (32) and (36) :

$$\begin{aligned} \underline{\nabla} \cdot \underline{E} &= -\nabla^2 \phi + \underline{\nabla} \cdot (\phi \underline{\omega}) - (54) \\ &= -\underline{\omega} \cdot \underline{\nabla} \phi + \phi \omega^2 = \frac{f}{\epsilon} \end{aligned}$$

3) The charge density is :

$$\rho = \epsilon_0 \left( \phi \omega^2 - \omega \frac{\partial \phi}{\partial z} \right) \quad - (55)$$

where

$$\phi = - \frac{e}{4\pi\epsilon_0} \frac{1}{z} \quad - (56)$$

and

$$\omega = - \frac{2}{z} \quad - (57)$$

$$\begin{aligned} \text{So } \rho &= \frac{\epsilon_0 e}{4\pi\epsilon_0} \left( -\frac{4}{z^3} + \frac{2}{z^3} \right) \quad - (57) \\ &= - \frac{2e}{4\pi z^3} \end{aligned}$$

The volume occupied by the charge is :

$$V = 4\pi z^3 \quad - (58)$$

As in previous work per unit volume resonance energy from eq. (54) where charge is oscillating  
A finite charge density must be defined as

follows:

$$\rho_0 = - \frac{2e}{4\pi z_0^3} \quad - (59)$$

in a finite volume

$$V_0 = 4\pi z_0^3 \quad - (60)$$



9) The correct way of doing this integral is to use the divergence theorem:

$$\oint_S \underline{d} \cdot \underline{n} dA = \int_V \underline{\nabla} \cdot \underline{d} d^3x \quad - (61)$$

where  $\underline{d}(x)$  is any well behaved vector field in a volume  $V$  surrounded by the closed surface  $S$ .

Here  $\underline{n}$  is the outwardly directed normal.

The ECE Coulomb law is:

$$\underline{\nabla} \cdot \underline{E} = \underline{\omega} \cdot \underline{E} = \frac{\rho}{\epsilon_0} \quad - (62)$$

Its integral form is:

$$\begin{aligned} \oint_S \underline{E} \cdot \underline{n} dA &= 4\pi \int_V \rho(x) d^3x \\ &= 4\pi\epsilon_0 \int_V \underline{\omega} \cdot \underline{E} d^3x \end{aligned} \quad - (63)$$

i.e.

$$\boxed{\oint_S \underline{E} \cdot \underline{n} dA = 4\pi\epsilon_0 \int_V \underline{\omega} \cdot \underline{E} d^3x} \quad - (64)$$

These equations are satisfied by:

$$\boxed{\underline{n} = \frac{\underline{\omega}}{\omega}} \quad - (65)$$

<sup>(10)</sup> from which it follows automatically that

$$\underline{\nabla} \cdot \underline{E} = \underline{\omega} \cdot \underline{E} \quad - (66)$$

Q.E.D.