

74(7) : Cartesian Relations in the Beta Ellipse

The Cartesian representation of the Beta ellipse is:

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1 \quad - (1)$$

where

$$x = ae + r \cos \beta \quad - (2)$$

$$y = r \sin \beta \quad - (3)$$

and where a and b are the semimajor and semiminor axes. Eqs. (1) to (3) correspond to:

$$r = \frac{d}{1 + e \cos \beta} \quad - (4)$$

$$d = a(1 - e^2) \quad - (5)$$

$$e^2 = 1 - \frac{b^2}{a^2} \quad - (6)$$

and to the Hamiltonian:

$$H = \frac{1}{2}mv^2 - \frac{k}{r} \quad - (7)$$

where

$$v^2 = \dot{r}^2 + \dot{\beta}^2 r^2 \quad - (8)$$

$$\dot{\beta}^2 = \dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2 \quad - (9)$$

and the Lagrangian:

$$\mathcal{L} = \frac{1}{2}mv^2 + \frac{k}{r} \quad - (10)$$

The possible Euler Lagrange equations are:

$$2) \quad \frac{\partial \mathcal{L}}{\partial r} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{r}} \right) \quad - (11)$$

$$\frac{\partial \mathcal{L}}{\partial \beta} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\beta}} \right) \quad - (12)$$

$$\frac{\partial \mathcal{L}}{\partial \theta} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) \quad - (13)$$

$$\text{and} \quad \frac{\partial \mathcal{L}}{\partial \phi} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \quad - (14)$$

Eqs. (10) to (14) give:

$$m\ddot{r} = mr\dot{\beta}^2 - \frac{L^2}{mr^3} \quad - (15)$$

$$\frac{d}{dt} (mr^2\dot{\beta}) = 0 \quad - (16)$$

$$\frac{d}{dt} (mr^2 \sin^2 \theta \dot{\phi}) = 0 \quad - (17)$$

$$\frac{d}{dt} (mr^2 \dot{\theta}) = 2 \sin \theta \cos \theta \dot{\phi}^2 \quad - (18)$$

From basic considerations as in UFT 269:

$$L_z = mr^2 \sin^2 \theta \dot{\phi} \quad - (19)$$

$$L^2 = m^2 r^4 (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) \quad - (20)$$

$$\text{so eq. (16) is:} \quad \frac{dL}{dt} = 0 \quad - (21)$$

3) and eq. (17) is:

$$\frac{dL_z}{dt} = 0 \quad - (22)$$

Eq. (18) shows that $mr^2\dot{\theta}$ is not a constant of motion, but eqs. (16) and (17) mean that L

and L_z are constants of motion.

From eqs. (19) and (20):

$$\frac{d\theta}{dt} = \frac{1}{mr^2} \left(L^2 - \frac{L_z^2}{\sin^2 \theta} \right)^{1/2} \quad - (23)$$

and

$$\frac{d\beta}{dt} = \frac{L}{mr^2} \quad - (24)$$

$$\frac{d\phi}{dt} = \frac{L_z}{mr^2 \sin^2 \theta} \quad - (25)$$

Eqs. (23) to (25) are true for a kinetic

energy

$$T = \frac{1}{2} m v^2 = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\beta}^2)$$

$$= \frac{1}{2} m \left(\dot{r}^2 + r^2 (\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2) \right) \quad - (26)$$

and for any potential energy U .

It follows that:

$$4) \quad \frac{d\beta}{d\theta} = \frac{L}{\left(\frac{L^2 - L_z^2}{\sin^2 \theta} \right)^{1/2}} \quad - (27)$$

$$\text{and } \frac{d\beta}{d\phi} = \left(\frac{L}{L_z} \right) \sin^2 \theta \quad - (28)$$

Eqs. (27) and (28) give the relation between β and θ and β and ϕ .

Therefore:

$$\beta = \int \frac{L d\theta}{\left(\frac{L^2 - L_z^2}{\sin^2 \theta} \right)^{1/2}} = -\sin^{-1} \left(\frac{L \cos \theta}{\left(L^2 - L_z^2 \right)^{1/2}} \right) \quad - (29)$$

$$\text{and } \sin \beta = - \frac{L \cos \theta}{\left(L^2 - L_z^2 \right)^{1/2}} \quad - (30)$$

It follows that:

$$\begin{aligned} \cos^2 \beta &= 1 - \sin^2 \beta \\ &= 1 - \frac{L^2 \cos^2 \theta}{\left(L^2 - L_z^2 \right)} \end{aligned} \quad - (31)$$

From Eq. (28):

$$\beta = \int \frac{L}{L_z} \sin^2 \theta d\phi \quad - (32)$$

and

$$\phi = \frac{L_z}{L} \int \frac{d\beta}{\sin^2 \theta} \quad - (33)$$

From Eq. (30):

$$\cos^2 \theta = \sin^2 \beta \quad - (34)$$

$$\left(\frac{L^2}{L^2 - L_z^2} \right) \cos^2 \theta = \left(1 - \left(\frac{L_z}{L} \right)^2 \right) \sin^2 \beta \quad - (35)$$

So

$$\sin^2 \theta = 1 - \cos^2 \theta = 1 - \left(1 - \left(\frac{L_z}{L} \right)^2 \right) \sin^2 \beta \quad - (36)$$

Therefore:

$$\phi = \frac{L_z}{L} \int \left(1 - \left(1 - \left(\frac{L_z}{L} \right)^2 \right) \sin^2 \beta \right)^{-1} d\beta \quad - (37)$$

$$= \tan^{-1} \left(\frac{L_z}{L} \tan \beta \right)$$

Therefore:

$$\tan \phi = \frac{L_z}{L} \tan \beta \quad - (38)$$

It follows that:

$$\tan \beta = \frac{\sin \beta}{\cos \beta} = \frac{(1 - \cos^2 \beta)^{1/2}}{\cos \beta} = \frac{L}{L_2} \tan \phi \quad - (39)$$

so:

$$1 - \cos^2 \beta = \frac{L^2}{L_2^2} \tan^2 \phi \cos^2 \beta \quad - (40)$$

and

$$\begin{aligned} \cos^2 \beta &= \frac{1}{1 + \left(\frac{L}{L_2}\right)^2 \tan^2 \phi} \quad - (41) \\ &= \frac{\cos^2 \phi}{\sin^2 \phi + \left(\frac{L}{L_2}\right)^2 \cos^2 \phi} \end{aligned}$$

Therefore:

$$X = a_e + \frac{r \cos \phi}{\left(\sin^2 \phi + \left(\frac{L}{L_2}\right)^2 \cos^2 \phi\right)^{1/2}} \quad - (42)$$

$$\begin{aligned} Y &= \frac{r \cos \theta}{\left(1 - \left(\frac{L_2}{L}\right)^2\right)^{1/2}} = \frac{Z}{\left(1 - \left(\frac{L_2}{L}\right)^2\right)^{1/2}} \quad - (43) \end{aligned}$$

7) Eq. (43) shows clearly that:

$$Z = \left(1 - \left(\frac{L_z}{L} \right)^2 \right)^{1/2} Y \quad - (44)$$

$\rightarrow 0$

as $L \rightarrow L_z$ - (45)

This is an indication that the orbit is three dimensional.

The equation linking Z and X is found from eqs. (1) and (44):

$$\frac{X^2}{a^2} + \frac{Z^2}{b^2 \left(1 - \left(\frac{L_z}{L} \right)^2 \right)} = 1 \quad - (46)$$

where

$$\frac{X^2}{a^2} + \frac{Y^2}{b^2} = 1 \quad - (47)$$
