

### 05(3) : Gravitationally affected wave functions of the H Atom

As in the previous note the Schrodinger equation in presence of the earth's gravitational field is:

$$\left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r} + V_1 \right) \psi = E \psi \quad (1)$$

where

$$V_1 = -m_2 \frac{MG}{R} \quad (2)$$

So:

$$\langle E \rangle = \int \psi^* E \psi d\tau$$

$$= \int \psi^* \left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r} + V_1 \right) \psi d\tau \quad (3)$$

$$= - \left( \frac{m e^4}{32\pi^2 \epsilon_0^2 \hbar^2} \right) \frac{1}{n^2} - \frac{m_2 MG}{R}$$

This means that the absolute values of the energy levels are shifted by  $V_1$ . The atomic spectrum is not affected because it depends on transitions between different  $n$ .

However eq. (1) provides additional theoretical information about the H atom in the gravitational field as follows.

First write the potential term in eq. (1) in terms of a coordinate  $r_1$  defined as:

$$2) \quad \frac{e^2}{4\pi\epsilon_0 r_1} = \frac{e^2}{4\pi\epsilon_0 r} + V_1 \quad - (4)$$

So: 
$$\frac{1}{r_1} = \frac{1}{r} + x \quad - (5)$$

where 
$$x = \frac{4\pi\epsilon_0 m_2 M G}{e^2 R} \quad - (6)$$

So: 
$$\boxed{r_1 = \frac{r}{1 + x r}} \quad - (7)$$

From eq. (6): 
$$x = 4.5346 \times 10^8 \text{ m}^{-1} \quad - (8)$$

Now consider the Schrodinger equation in the coordinate system  $(r_1, \theta, \phi)$ . It is

$$\left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r_1^2} - \frac{e^2}{4\pi\epsilon_0 r_1} \right) \psi = E_1 \psi \quad - (9)$$

for which 
$$E_1 = - \left( \frac{m e^4}{32\pi^2 \epsilon_0^2 \hbar^2} \right) \frac{1}{n^2} \quad - (10)$$

The use of the coordinate system  $(r_1, \theta, \phi)$  has the effect of giving the hydrogenic energy levels (10) in the presence of the earth's gravitational field. The wave functions of eq. (9) are:

$$3) \psi(r, \theta, \phi) = R(r) Y(\theta, \phi) \quad - (11)$$

where:

$$R_{nl}(r) = \left( \frac{2Z}{na} \right) \left[ \frac{(n-l-1)!}{2n[(n+l)!]^3} \right] \rho_1^l L_{n+l}^{2l+1}(\rho_1) \exp\left(-\frac{\rho_1}{2}\right) \quad - (12)$$

where:

$$\rho_1 = \left( \frac{2Z}{na} \right) r \quad - (13)$$

Here  $a$  is the Bohr radius and  $Z$  the atomic number.  
The expectation value is defined by:

$$\begin{aligned} \langle E_1 \rangle &= \int \psi_1^* E_1 \psi_1 d\tau \quad - (14) \\ &= \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^{\infty} \psi_1^* E_1 \psi_1 r^2 \sin\theta dr d\theta d\phi \end{aligned}$$

Therefore  $r$  of the usual Schrodinger equation:

$$\left( -\frac{\hbar^2}{2m} \nabla^2 - \frac{e^2}{4\pi\epsilon_0 r} \right) \psi = E \psi \quad - (15)$$

is replaced by  $r_1$  defined from eq. (1).  
This procedure gives the same energy levels (10) but for eq. (9) and eq. (15), but the wave functions of eq. (9) are different. The wave functions in the coordinate system  $(r, \theta, \phi)$  can be expressed in terms of  $r$  using eq. (7). The first few radial wave functions

4) of the coordinate system  $(r, \theta, \phi)$  are as follows.

$$\underline{1s \text{ Orbital}}, \underline{n=1}, \underline{l=0}$$

$$R_{10}(r) = \left(\frac{1}{a}\right)^{3/2} \exp\left(-\frac{\rho_1}{2}\right) \quad (16)$$

where

$$\rho_1 = \frac{2Z}{na} \left(\frac{r}{1+\alpha r}\right) \quad (17)$$

$$\underline{2s \text{ Orbital}}, \underline{n=2}, \underline{l=0}$$

$$R_{20}(r) = \left(\frac{1}{a}\right)^{3/2} \left(\frac{1}{2\sqrt{2}}\right) (2-\rho_1) \exp\left(-\frac{\rho_1}{2}\right) \quad (18)$$

$$\underline{2p \text{ Orbital}}, \underline{n=2}, \underline{l=1}$$

$$R_{21}(r) = \left(\frac{1}{a}\right)^{3/2} \left(\frac{1}{2\sqrt{6}}\right) \rho_1 \exp\left(-\frac{\rho_1}{2}\right) \quad (19)$$

$$\underline{3s \text{ Orbital}}, \underline{n=3}, \underline{l=0}$$

$$R_{30}(r) = \left(\frac{1}{a}\right)^{3/2} \left(\frac{1}{9\sqrt{3}}\right) (6-6\rho_1+\rho_1^2) \exp\left(-\frac{\rho_1}{2}\right) \quad (20)$$

$$\underline{3p \text{ Orbital}}, \underline{n=3}, \underline{l=1}$$

$$R_{31}(r) = \left(\frac{1}{a}\right)^{3/2} \left(\frac{1}{9\sqrt{6}}\right) (4-\rho_1) \rho_1 \exp\left(-\frac{\rho_1}{2}\right) \quad (21)$$

$$\underline{3d \text{ Orbital}}, \underline{n=3}, \underline{l=2}$$

$$R_{32}(r) = \left(\frac{1}{a}\right)^{3/2} \left(\frac{1}{9\sqrt{30}}\right) \rho_1^2 \exp\left(-\frac{\rho_1}{2}\right) \quad (22)$$

5) The complete wave functions are given in each of eqs. (16) to (22) by:

$$\begin{aligned}\psi_1 &= R(r_1) Y(\theta, \phi) \quad - (23) \\ &= R\left(\frac{r}{1+xr}\right) Y(\theta, \phi)\end{aligned}$$

Eqs. (16) to (22) can be plotted by computer and give entirely new information about the H atom in the gravitational field of a system of mass M. This is true for all atoms and molecules.

The energy expectation values are:

$$\langle E_1 \rangle = \int \psi_1 E_1 \psi_1^* d\tau_1 \quad - (24)$$

$$= - \left( \frac{ne^4}{32\pi^2 \epsilon_0 \hbar^2} \right) \frac{1}{n^2} \quad - (25)$$

where:

$$\left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r_1^2} - \frac{e^2}{4\pi \epsilon_0 r_1} \right) \psi_1 = E_1 \psi_1$$

These are the same as the energy expectation values of

$$\left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} - \frac{e^2}{4\pi \epsilon_0 r} \right) \psi = E \psi \quad - (26)$$