

311(7): (checking the Inhomogeneous Field Equations

these are derived from the Cartan-Evans identity:

$$D_\mu T^{a\mu\nu} = R^a{}_\mu{}^{\mu\nu} - (1)$$

i.e.

$$\partial_\mu T^{a\mu\nu} = R^a{}_\mu{}^{\mu\nu} - \omega^a{}_{\mu b} T^{b\mu\nu} - (2)$$

using for each  $a$ :

$$T^{\mu\nu} = \begin{bmatrix} 0 & -T^1(ab) & -T^2(ab) & -T^3(ab) \\ T^1(ab) & 0 & -T^3(sp) & T^2(sp) \\ T^2(ab) & T^3(sp) & 0 & -T^1(sp) \\ T^3(ab) & -T^2(sp) & T^1(sp) & 0 \end{bmatrix} - (3)$$

The Coulomb law is derived from:

$$\nu = 0 - (4)$$

i.e.

$$\begin{aligned} & \partial_1 T^{a10} + \partial_2 T^{a20} + \partial_3 T^{a30} - (5) \\ & = R^a{}_{110} + R^a{}_{220} + R^a{}_{330} - \omega^a{}_{1b} T^{b10} - \omega^a{}_{2b} T^{b20} - \omega^a{}_{3b} T^{b30} \end{aligned}$$

$$\underline{\nabla} \cdot \underline{T}^a(ab) = \underline{\omega}^a{}_b \cdot \underline{T}^b(ab) - \underline{q}^b \cdot \underline{R}^a{}_b(ab) - (6)$$

where:

$$\underline{T}^a(ab) = T^{a1}(ab) \underline{i} + T^{a2}(ab) \underline{j} + T^{a3}(ab) \underline{k}$$

$$\text{and } \underline{R}^a{}_b(ab) = R^{a1}{}_b(ab) \underline{i} + R^{a2}{}_b(ab) \underline{j} + R^{a3}{}_b(ab) \underline{k} - (7)$$

-(8)

2) Now use:

$$\underline{E}^a = c A^{(0)} \underline{T}^a(\alpha) - (9)$$

and

$$\underline{A}^a = A^{(0)} \underline{v}^a - (10)$$

to find that:

$$\underline{\nabla} \cdot \underline{E}^a = \underline{\omega}^a_b \cdot \underline{E}^b - c A^b \cdot \underline{R}^a_b(\alpha) - (11)$$

This is the same as UFT255 eq. (66) and Slide 25 of  
 @ Engineering Model, Q.E.D.

Now use:

$$\underline{E}^a_b = c W^{(0)} \underline{R}^a_b(\alpha) - (11)$$

to find that:

$$\underline{\nabla} \cdot \underline{E}^a = \underline{\omega}^a_b \cdot \underline{E}^b - \frac{1}{W^{(0)}} A^b \cdot \underline{E}^a_b - (12)$$

$$= \underline{\omega}^a_b \cdot \underline{E}^b - \frac{1}{r^{(0)}} \underline{v}^b \cdot \underline{E}^a_b$$

Now remove indices using the methods developed in  
 preceding notes and pages to find that:

$$\underline{\nabla} \cdot \underline{E} = 2 \underline{E} \cdot \left( \frac{1}{r^{(0)}} \underline{v} - \underline{\omega} \right) - (13)$$

3) This is the same as in notes 317(2) and 317(4), QED.

Therefore:

$$\boxed{\begin{aligned}\underline{\nabla} \cdot \underline{E} &= 2 \underline{E} \cdot \underline{\kappa} \\ \underline{\nabla} \cdot \underline{B} &= 2 \underline{B} \cdot \underline{\kappa}\end{aligned}} \quad - (14)$$

where:  $\underline{\kappa} = \frac{1}{r(\omega)} \underline{\nabla} - \underline{\omega} \quad - (15)$

In ECE2 there is exact symmetry between the Gauss law of magnetism and the Coulomb law.

The Ampère Maxwell law is obtained from:

$$\sim = 1, 2, 3 \quad - (16)$$

in eq. (2). For:  $\sim = 1 \quad - (17)$

it follows that:

$$\begin{aligned}& \partial_0 T^{a01} + \partial_2 T^{a21} + \partial_3 T^{a31} \\ &= \gamma^b_0 R^a_b{}^{01} + \gamma^b_2 R^a_b{}^{21} + \gamma^b_3 R^a_b{}^{31} \\ &\quad - \omega^a_{0b} T^{b01} - \omega^a_{2b} T^{b21} - \omega^a_{3b} T^{b31}\end{aligned} \quad - (18)$$

$$\begin{aligned}\text{i.e. } & -\partial_0 T^1(\omega) + \partial_2 T^3(s_p) - \partial_3 T^2(s_p) \\ &= -\gamma^b_0 R^a_b{}^1(\omega) + \gamma^b_2 R^a_b{}^3(s_p) - \gamma^b_3 R^a_b{}^2(s_p) \\ &\quad + \omega^a_{0b} T^{b1}(\omega) - \omega^a_{2b} T^{b3}(s_p) + \omega^a_{3b} T^{b2}(s_p)\end{aligned} \quad - (19)$$

4) This means that:

$$\begin{aligned}
 & -\frac{1}{c} \frac{dT_x(orb)}{dt} + (\underline{\nabla} \times \underline{T}(sp))_x \\
 & = \omega^a_{ob} T^b_x(orb) + (\underline{\omega}^a_b \times \underline{T}^b(sp))_x \\
 & \quad - \underline{v}^b_o R^a_{bx}(orb) - (\underline{v}^b \times \underline{R}^a_b(sp))_x
 \end{aligned} \tag{20}$$

There are similar results for the y and z components

So:

$$\begin{aligned}
 & \underline{\nabla} \times \underline{T}^a(sp) - \frac{1}{c} \frac{dT^a(orb)}{dt} \\
 & = \omega^a_{ob} \underline{T}^b(orb) - \underline{v}^b_o \underline{R}^a_b(orb) \\
 & \quad + \underline{\omega}^a_b \times \underline{T}^b(sp) - \underline{v}^b \times \underline{R}^a_b(sp)
 \end{aligned} \tag{21}$$

which is eq. (21) of HFT255, QED.

Now use:  $\underline{B}^a = A^{(0)} \underline{T}^a(sp) \tag{22}$

$$\underline{E}^a = c A^{(0)} \underline{T}^a(orb) \tag{23}$$

$$\underline{B}^a_b = W^{(0)} \underline{R}^a_b(sp) \tag{24}$$

$$\underline{E}^a_b = c W^{(0)} \underline{R}^a_b(orb) \tag{25}$$

To find that:

5)

$$\underline{\nabla} \times \underline{B}^a - \frac{1}{c^2} \frac{\partial \underline{E}^a}{\partial t}$$

$$= A^{(0)} \left( \omega^a_{ob} \underline{T}^b(orb) - \underline{v}^b_o R^a_b(orb) \right) + \omega^a_b \times \underline{T}^b(sp) - \underline{v}^b \times R^a_b(sp)$$

$$= \frac{\omega^a_{ob}}{c} \underline{E}^b - A^b_o R^a_b(orb) + \omega^a_b \times \underline{B}^b - \underline{A}^b \times R^a_b(sp)$$

This is the same as in the Engineering Model,  
slide 25, Q.E.D. (26)

Now use eqs. (24) and (25) to find:

$$\underline{\nabla} \times \underline{B}^a - \frac{1}{c^2} \frac{\partial \underline{E}^a}{\partial t}$$

-(27)

$$= \frac{\omega^a_{ob}}{c} \underline{E}^b - \frac{A^b_o}{c W^{(0)}} \underline{E}^a_b + \omega^a_b \times \underline{B}^b - \frac{A^b \times \underline{B}^a_b}{W^{(0)}}$$

and remove indices to find that:

-(28)

$$\underline{\nabla} \times \underline{B} - \frac{1}{c^2} \frac{\partial \underline{E}}{\partial t} = 2 \left[ \left( \underline{g} - \frac{\underline{\omega}}{r^{(0)}} \right) \times \underline{B} - \frac{1}{c} \left( \underline{\omega}_o - \frac{\underline{v}_o}{r^{(0)}} \right) \underline{E} \right]$$

which is Note 317(b), Eq. (b), Q.E.D.

6) Final Format of Field Equations

$$\underline{\nabla} \cdot \underline{B} = 2 \underline{B} \cdot \underline{\kappa} \quad - (29)$$

$$\frac{\partial \underline{B}}{\partial t} + \underline{\nabla} \times \underline{E} = 2 \left( c \underline{\kappa}_0 \underline{B} + \underline{\kappa} \times \underline{E} \right) \quad - (30)$$

$$\underline{\nabla} \cdot \underline{E} = 2 \underline{E} \cdot \underline{\kappa} \quad - (31)$$

$$\underline{\nabla} \times \underline{B} - \frac{1}{c^2} \frac{\partial \underline{E}}{\partial t} = 2 \left( -\underline{\kappa} \times \underline{B} - \frac{\underline{\kappa}_0}{c} \underline{E} \right) \quad - (32)$$

where:

$$\underline{\kappa}_0 = \underline{\omega}_0 - \frac{q \underline{v}_0}{r^{(0)}} \quad - (33)$$

and

$$\underline{\kappa} = \underline{\omega} - \frac{q \underline{v}}{r^{(0)}} \quad - (34)$$

There is a symmetry between the  
homogeneous and inhomogeneous field equations.

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