

40 (6): Cyclic Symmetry and B (3) Field from the  
W Potentials

Using the result:  $\underline{W}^a = \underline{A}^a$  — (1)

it follows that:

$$\underline{B}^{(1)*} = \underline{\nabla} \times \underline{W}^{(1)*} - \frac{i}{\underline{W}^{(0)}} \underline{W}^{(2)} \times \underline{W}^{(3)} - (2)$$

$$\underline{B}^{(2)*} = \underline{\nabla} \times \underline{W}^{(2)*} - \frac{i}{\underline{W}^{(0)}} \underline{W}^{(3)} \times \underline{W}^{(1)} - (3)$$

$$\underline{B}^{(3)*} = \underline{\nabla} \times \underline{W}^{(3)*} - \frac{i}{\underline{W}^{(0)}} \underline{W}^{(1)} \times \underline{W}^{(2)} - (4)$$

in the complex circular basis:

$$\underline{e}^{(1)} \times \underline{e}^{(2)} = i \underline{e}^{(3)*} - (5)$$

et cetera

The Aharonov Bohm vacuum is therefore defined

by:  $\underline{\nabla} \times \underline{W}^{(1)*} = \frac{i}{\underline{W}^{(0)}} \underline{W}^{(2)} \times \underline{W}^{(3)} - (6)$

et cetera

i.e. by a cyclic relation between spin connection vectors:

$$\underline{\nabla} \times \underline{\omega}^{(1)*} = i \underline{\omega}^{(2)} \times \underline{\omega}^{(3)} - (7)$$

et cetera

Under this condition there is a  $\underline{W}^a$  potential

) In magnetic field.

From Cartan geometry the magnetic field is also defined by:

$$\underline{B}^a = \underline{\nabla} \times \underline{W}^a - \underline{\omega}^a{}_b \times \underline{W}^b \quad (8)$$

This equation must be used in the complex circular basis, where cross products must be defined as in Eq. (5). For example, the index of the  $\underline{B}^{(3)}$  field is:

$$a = (3) - (9)$$

So:

$$\begin{aligned} \underline{B}^{(3)*} &= \underline{B}^{(3)} = \underline{\nabla} \times \underline{W}^{(3)*} - i \underline{\omega}^{(2)} \times \underline{W}^{(2)} \\ &= \underline{\nabla} \times \underline{W}^{(3)*} - i \underline{\omega}^{(1)} \times \underline{W}^{(2)} \quad (10) \end{aligned}$$

so

$$\boxed{\underline{\omega}^{(3)*}_{(2)} = \underline{\omega}^{(1)}} \quad (11)$$

In general there is summation over repeated b indices in Eq. (8), and in the complex circular basis:

$$\begin{aligned} \underline{B}^{(3)*} &= \underline{\nabla} \times \underline{W}^{(3)*} - i \underline{\omega}^{(3)*}_{(1)} \times \underline{W}^{(1)} - (12) \\ &\quad - i \underline{\omega}^{(3)*}_{(2)} \times \underline{W}^{(2)} - i \underline{\omega}^{(3)*}_{(3)} \times \underline{W}^{(3)} \end{aligned}$$

but the basis (5) implies that the cross products  $-i \underline{\omega}^{(3)*}_{(1)} \times \underline{W}^{(1)}$  and  $-i \underline{\omega}^{(3)*}_{(3)} \times \underline{W}^{(3)}$  cannot give a vector in the direction of  $\underline{B}^{(3)*}$ . This implies the result (11).

Similarly:

$$\underline{\omega}^{(2)} = \underline{\omega}^{(1)*} \quad - (13)$$

$$\underline{\omega}^{(1)} = \underline{\omega}^{(3)*} \quad - (14)$$

This implies:

$$\underline{\omega}^{(3)*} = \epsilon^{(3)*}_{(1)(2)} \underline{\omega}^{(1)} \quad - (15)$$

$$\underline{\omega}^{(1)*} = \epsilon^{(1)*}_{(2)(3)} \underline{\omega}^{(2)} \quad - (16)$$

$$\underline{\omega}^{(2)*} = \epsilon^{(2)*}_{(3)(1)} \underline{\omega}^{(3)} \quad - (17)$$

where  $\epsilon^{(a)*}_{(b)(c)}$  is a totally antisymmetric unit tensor in three dimensions and in the complex circular basis. Under these conditions eq. (8) becomes:

$$\underline{B}^{(1)*} = \underline{\nabla} \times \underline{W}^{(1)*} - i \underline{\omega}^{(2)} \times \underline{W}^{(3)} \quad - (18)$$

& cyclicum, which is consistent with:

$$\underline{W}^a = \underline{W}^{(a)} \underline{\omega}^a \quad - (19)$$

Q.E.D.

In general:

$$\underline{W}^{(a)} = \underline{W}^{(a)} \underline{\omega}^{(a)}, \quad - (20)$$

$$a = 1, 2, 3$$

$$\underline{W}^{(1)} \times \underline{W}^{(2)} = i \underline{W}^{(1)} \underline{W}^{(3)*} \quad - (21)$$

which is the W cyclic theorem.

4) If it is assumed that  $\underline{W}^{(1)}$  and  $\underline{W}^{(2)}$  are plane waves:

$$\underline{W}^{(1)} = \underline{W}^{(2)*} = \frac{\underline{W}^{(0)}}{\sqrt{5}} (\underline{i} - i\underline{j}) e^{i\phi} \quad - (20)$$

then:

$$\underline{W}^{(3)} = \underline{W}^{(0)} \underline{k} \quad - (21)$$

and

$$\underline{\nabla} \times \underline{W}^{(3)} = \underline{0} \quad - (22)$$

so:

$$\underline{B}^{(3)} = \underline{B}^{(3)*} = -\frac{i}{\underline{W}^{(0)}} \underline{W}^{(1)} \times \underline{W}^{(2)} \quad - (23)$$

This assume that the spin connection vectors are

plane waves:

$$\underline{\omega}^{(1)} = \underline{\omega}^{(2)*} = \frac{\underline{\omega}^{(0)}}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} \quad - (24)$$


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