

541(3): The Gravitomagnetic Minimal Prescription.

Consider the minimal prescription in electrodynamics:

$$\underline{p} \rightarrow \underline{p} - e \underline{A} \quad - (1)$$

In ECE2 theory this goes to:

$$\underline{p} \rightarrow \underline{p} - e \underline{W} \quad - (2)$$

where $-e$ is the charge of the electron. In the presence of \underline{W} the Hamiltonian of a free particle becomes:

$$\begin{aligned} H &= \frac{1}{2m} (\underline{p} - e \underline{W}) \cdot (\underline{p} - e \underline{W}) \\ &= \frac{\underline{p}^2}{2m} + \frac{e^2 \underline{W}^2}{2m} - \frac{e}{m} \underline{p} \cdot \underline{W} \quad - (3) \end{aligned}$$

The magnetic field is defined by:

$$\underline{B} = \nabla \times \underline{W} \quad - (4)$$

and for a uniform magnetic field:

$$\underline{W} = \frac{1}{2} \underline{B} \times \underline{r} \quad - (5)$$

Therefore:

$$\begin{aligned} H_1 &= -\frac{e}{m} \underline{p} \cdot \underline{W} = -\frac{e}{m} \underline{p} \cdot \underline{B} \times \underline{r} \\ &= -\frac{e}{2m} \underline{L} \cdot \underline{B} \quad - (6) \end{aligned}$$

where

$$\underline{L} = \underline{r} \times \underline{p} \quad - (7)$$

the orbital angular momentum. The magnetic dipole moment is defined as:

$$\underline{m} = \frac{e}{2m} \underline{L} \quad - (8)$$

$$H_1 = -\underline{m} \cdot \underline{B} \quad - (9)$$

2) The torque generated between \underline{m} and \underline{B} is:

$$\underline{\tau}_g = \underline{m} \times \underline{B} \quad - (10)$$

and the Larmor precession frequency is:

$$\omega = \frac{eB}{2m} \quad - (11)$$

as observed in electron spin resonance (ESR).

This well known theory can be developed for any astrominial precession as follows.
The precession is due to the gravitomagnetic field:

$$\underline{\Omega}_g = \underline{\nabla} \times \underline{v}_g \quad - (12)$$

where

$$\underline{W}_g = \underline{v}_g \quad - (13)$$

Therefore the gravitomagnetic field is a vorticity in ECE2 spacetime.

The gravitomagnetic minimal prescription is:

$$\underline{p} \rightarrow \underline{p} + m \underline{v}_g \quad - (14)$$

for a free particle \underline{p} is the preserve of the potential defined in Eq. (13). The free particle Hamiltonian becomes:

$$H = \frac{1}{2m} (\underline{p} + m \underline{v}_g) \cdot (\underline{p} + m \underline{v}_g) \quad - (15)$$

where:

$$\underline{v}_g = \frac{1}{2} \underline{\Omega}_g \times \underline{r} \quad - (16)$$

for a uniform gravitomagnetic field $\underline{\Omega}_g$. Therefore the free particle Hamiltonian becomes:

$$H = \frac{\underline{p}^2}{2m} + \frac{1}{2} m \underline{v}_g^2 + \frac{1}{2} \underline{L} \cdot \underline{\Omega}_g \quad - (17)$$

where \underline{L} is defined by eq. (7). The factor $1/2$ in the last term of eq. (17) takes the place of $-e/2m$ in leishmanics, so the gyromagnetic Larmor precession frequency is:

$$\boxed{\Omega = \frac{1}{2} \Omega_g} \quad - (18)$$

In the presence of a precession Ω is radians per second the orbital function is changed from:

$$H = \frac{p^2}{2m} + U \quad - (19)$$

to

$$\boxed{H = \frac{p^2}{2m} + \frac{1}{2} m v_g^2 + \Omega L + U} \quad - (20)$$

and the Lagrangian becomes:

$$\boxed{L = \frac{p^2}{2m} + \frac{1}{2} m v_g^2 + \Omega L - U} \quad - (21)$$

The orbital characteristics can now be worked out by development of standard methods, to give the precessing orbit.