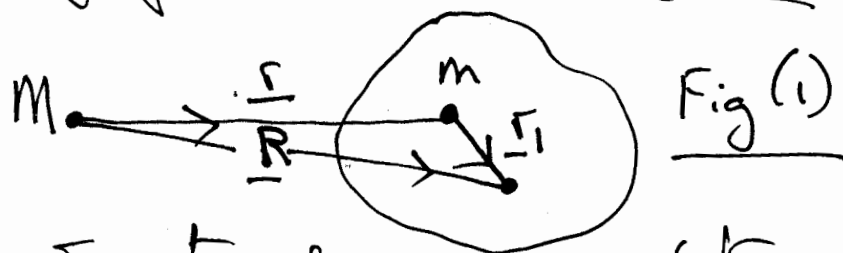


# 369(a): General Theory of Milankovitch Cycles

Fig (1)



Consider any asymmetric top of mass  $m$  orbiting the sun of mass  $M$ . The distance between the centre of mass of the object and the sun is  $\underline{r}$ . The distance of any point in the object from its centre of mass is  $\underline{r}_1$ . Therefore the distance between the sun and this point is  $\underline{R} = \underline{r} + \underline{r}_1$  - (1)

The Lagrangian of relevance is:

$$\mathcal{L} = \frac{1}{2} m (\dot{\underline{r}} + \dot{\underline{r}}_1) \cdot (\dot{\underline{r}} + \dot{\underline{r}}_1) - U \quad - (2)$$

where

$$U = -\frac{mMG}{r} \quad - (3)$$

The position of the centre of mass is denoted:

$$\underline{r} = x \underline{i} + y \underline{j} + z \underline{k} \quad - (4)$$

in the stationary frame.

The position of  $\underline{r}_1$  is:

$$\underline{r}_1 = r_{11} \underline{e}_1 + r_{12} \underline{e}_2 + r_{13} \underline{e}_3 \quad - (5)$$

in the frame (1, 2, 3) of the principal moments of inertia

2) of asymptotic top.

The transformation from  $(1, 2, 3)$  to  $(x, y, z)$  is given by:

$$\begin{bmatrix} \underline{x}_1 \\ \underline{x}_2 \\ \underline{x}_3 \end{bmatrix} = \begin{bmatrix} \cos\psi & \sin\psi & 0 \\ -\sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \cos\theta & 0 & \sin\theta \\ 0 & 1 & 0 \\ -\sin\theta & 0 & \cos\theta \end{bmatrix} \begin{bmatrix} \cos\phi & \sin\phi & 0 \\ -\sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \underline{i} \\ \underline{j} \\ \underline{k} \end{bmatrix} \quad - (6)$$

where  $\theta$ ,  $\phi$  and  $\psi$  are the Euler angles defined in Mariani and Thornton for example. So:

$$\begin{bmatrix} \underline{x}_1 \\ \underline{x}_2 \\ \underline{x}_3 \end{bmatrix} = \begin{bmatrix} A & B & C \\ D & E & F \\ G & H & I \end{bmatrix} \begin{bmatrix} \underline{i} \\ \underline{j} \\ \underline{k} \end{bmatrix} \quad - (7)$$

where:

$$A = \cos\psi \cos\phi - \cos\theta \sin\psi \sin\phi$$

$$B = -\sin\psi \cos\phi - \cos\theta \sin\psi \cos\phi$$

$$C = \sin\theta \cos\phi$$

$$D = \cos\psi \cos\phi + \cos\theta \sin\psi \sin\phi$$

$$E = -\sin\psi \sin\phi + \cos\theta \cos\psi \cos\phi$$

$$F = -\sin\theta \sin\phi$$

$$G = \sin\psi \cos\theta$$

$$H = \cos\psi \sin\theta$$

$$I = \cos\theta$$

- (8)

It follows that:

$$\begin{aligned}
 \underline{r}_1 &= r_{11}(\underline{A}\underline{i} + \underline{B}\underline{j} + \underline{C}\underline{k}) \\
 &+ r_{12}(\underline{D}\underline{i} + \underline{E}\underline{j} + \underline{F}\underline{k}) \\
 &+ r_{13}(\underline{G}\underline{i} + \underline{H}\underline{j} + \underline{I}\underline{k}) \\
 &= (r_{11}\underline{A} + r_{12}\underline{D} + r_{13}\underline{G})\underline{i} \\
 &+ (r_{11}\underline{B} + r_{12}\underline{E} + r_{13}\underline{H})\underline{j} \\
 &+ (r_{11}\underline{C} + r_{12}\underline{F} + r_{13}\underline{I})\underline{k} \\
 &= X_1\underline{i} + Y_1\underline{j} + Z_1\underline{k}
 \end{aligned} \tag{9}$$

where:

$$\begin{aligned}
 X_1 &= r_{11}\underline{A} + r_{12}\underline{D} + r_{13}\underline{G} \\
 Y_1 &= r_{11}\underline{B} + r_{12}\underline{E} + r_{13}\underline{H} \\
 Z_1 &= r_{11}\underline{C} + r_{12}\underline{F} + r_{13}\underline{I}
 \end{aligned} \tag{10}$$

The Lagrangian is therefore:

$$\begin{aligned}
 \mathcal{L} &= \frac{1}{2}m((\dot{x} + \dot{X}_1)^2 + (\dot{y} + \dot{Y}_1)^2 + (\dot{z} + \dot{Z}_1)^2) \\
 &+ \frac{mmG}{r} \tag{11}
 \end{aligned}$$

If the orbit is considered to be planar, then:

$$x = r \cos \theta_1 \tag{12}$$

$$y = r \sin \theta_1 \tag{13}$$

$$z = 0 \tag{14}$$

The system is governed by five Euler Lagrange equations in five Lagrange variables,  $r, \theta, \theta_1, \phi$  and  $\dot{\phi}$ :

$$\frac{\partial \mathcal{L}}{\partial r} = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{r}} \right) \quad - (15)$$

$$\frac{\partial \mathcal{L}}{\partial \theta_1} = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\theta}_1} \right) \quad - (16)$$

$$\frac{\partial \mathcal{L}}{\partial \theta} = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) \quad - (17)$$

$$\frac{\partial \mathcal{L}}{\partial \phi} = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \quad - (18)$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \ddot{\phi}} \right) \quad - (19)$$

These five equations must be solved simultaneously.

In eq. (11):

$$\dot{x} = \dot{r} \cos \theta_1 - r \dot{\theta}_1 \sin \theta_1 \quad - (20)$$

$$\dot{y} = \dot{r} \sin \theta_1 + r \dot{\theta}_1 \cos \theta_1 \quad - (21)$$

$$\dot{z} = 0 \quad - (22)$$

The distances  $r_{11}, r_{12}$  and  $r_{13}$  are constants related to the principal moments of inertia  $I_1, I_2$  and  $I_3$  of the asymmetric top.

) For the purposes of computing these can be used as  
 put parameters. Therefore:

$$\dot{X}_1 = r_{11} \frac{dA}{dt} + r_{12} \frac{dD}{dt} + r_{13} \frac{dB}{dt} - (23)$$

$$\dot{Y}_1 = r_{11} \frac{dB}{dt} + r_{12} \frac{dE}{dt} + r_{13} \frac{dH}{dt} - (24)$$

$$\dot{Z}_1 = r_{11} \frac{dC}{dt} + r_{12} \frac{dF}{dt} + r_{13} \frac{dI}{dt} - (25)$$

and can be evaluated by computer algebra to  
 eliminate human error.

The Lagrangian is therefore:

$$L = \frac{1}{2} m \left( \left( \dot{r} \cos \theta_1 - r \dot{\theta}_1 \sin \theta_1 + \dot{X}_1 \right)^2 \right. \\ \left. + \left( \dot{r} \sin \theta_1 + r \dot{\theta}_1 \cos \theta_1 + \dot{Y}_1 \right)^2 \right. \\ \left. + \dot{Z}_1^2 \right) + \frac{nmG}{r} - (26)$$

It is seen that the Manhewtel cycles are the  
 result of intricate dynamics - the motion of the  
 rotating asymmetric top superposed on its orbital  
 motion. All the motions are inter-related.

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