

394(6) . Vacuum Effects in the Magnetic Dipole Potential and Field

The magnetic dipole potential is :

$$\underline{A} = \frac{\mu_0}{4\pi} \frac{\underline{m} \times \underline{r}}{r^3} \quad - (1)$$

as in previous notation, and the magnetic field is :

$$\underline{B} = -\frac{\mu_0}{4\pi} \underline{m} \nabla^2 \left(\frac{1}{r} \right) + \frac{\mu_0}{4\pi r^3} \left(3 \underline{m} \cdot \frac{\underline{r} \underline{r}}{r^2} - \underline{m} \right)$$

$$\therefore = \underline{B}_c + \underline{B}_d \quad - (2)$$

in which the contact field is :

$$\begin{aligned} \underline{B}_c &= -\frac{\mu_0}{4\pi} \underline{m} \nabla^2 \left(\frac{1}{r} \right) \\ &= \mu_0 \underline{m} \delta_0(\underline{r}) \end{aligned} \quad - (3)$$

where $\delta_0(\underline{r})$ is the Dirac delta function.

The magnetic dipole field is :

$$\underline{B}_d = \frac{\mu_0}{4\pi r^3} \left(3 \underline{m} \cdot \frac{\underline{r} \underline{r}}{r^2} - \underline{m} \right) \quad - (4)$$

For each function in the presence of the vacuum we propose to postulate :

$$\underline{r} \rightarrow \underline{r} + \delta \underline{r} \quad - (5)$$

and

$$r = |\underline{r}| \rightarrow |\underline{r} + \delta \underline{r}| \quad - (6)$$

3) This postulate is applied to all equations of physics, saying that the r coordinate shivers in the presence of the vacuum.

As in previous work:

$$\begin{aligned} |\underline{r} + \delta \underline{r}| &= (r^2 + 2\underline{r} \cdot \delta \underline{r} + \delta \underline{r} \cdot \delta \underline{r})^{1/2} \\ &= r \left(1 + \frac{2\underline{r} \cdot \delta \underline{r} + \delta \underline{r} \cdot \delta \underline{r}}{r^2} \right)^{1/2} \\ &:= r(1+x)^{1/2} \quad - (7) \end{aligned}$$

where $x = (2\underline{r} \cdot \delta \underline{r} + \delta \underline{r} \cdot \delta \underline{r}) / r^2$ - (8)

So the magnetic dipole potential in the presence of the vacuum is:

$$\begin{aligned} \underline{A} &= \frac{\mu_0}{4\pi} \frac{\underline{m} \times (\underline{r} + \delta \underline{r})}{|\underline{r} + \delta \underline{r}|^3} \\ &= \frac{\mu_0}{4\pi r^3} \frac{\underline{m} \times (\underline{r} + \delta \underline{r})}{(1+x)^{3/2}} \quad - (9) \\ &= \frac{\mu_0 \underline{m} \times (\underline{r} + \delta \underline{r})}{4\pi r^3} (1+x)^{-3/2} \end{aligned}$$

Using the binomial expansion as developed by Newton:

$$(1+x)^{-3/2} = 1 - \frac{3x}{2} + \frac{15}{8}x^2 + \dots \quad - (10)$$

So in the presence of the vacuum:

$$\underline{A} = \frac{\mu_0}{4\pi r^3} \underline{m} \times (\underline{r} + \delta \underline{r}) \left(1 - \frac{3x}{2} + \frac{15}{8}x^2 + \dots \right) \quad - (11)$$

The average value $\langle A \rangle$ is finally evaluated using many terms as needed. The final result is graphed.

In note 393(4) it was shown that the dipole electric field in the presence of the vacuum is:

$$\underline{E} = \frac{1}{4\pi\epsilon_0} \left(\frac{(\underline{r} + \delta\underline{r})(\underline{p} \cdot (\underline{r} + \delta\underline{r}))}{r^5} \left(1 - \frac{5x}{2} + \frac{35}{8}x^2 + \dots \right) - \frac{\underline{p}}{r^3} \left(1 - \frac{3x}{2} + \frac{15x^2}{8} + \dots \right) \right) \quad (12)$$

The dipole magnetic field in the presence of the vacuum has the same structure:

$$\underline{B} = \frac{\mu_0}{4\pi} \left(\frac{(\underline{r} + \delta\underline{r})(\underline{p} \cdot (\underline{r} + \delta\underline{r}))}{r^5} \left(1 - \frac{5x}{2} + \frac{35}{8}x^2 + \dots \right) - \frac{\underline{p}}{r^3} \left(1 - \frac{3x}{2} + \frac{15x^2}{8} + \dots \right) \right) \quad (13)$$

Therefore after averaging it will display the same new features as the electric dipole field.
The magnet field will be worked out in the next note.