

15(8) : Lagrangian Method.

By definition:

$$\underline{r} = \frac{\underline{\dot{r}}}{m(r)^{1/2}} \underline{e}_r - (1)$$

and

$$\underline{\dot{r}} = \frac{1}{m(r)^{1/2}} (\dot{r} \underline{e}_r + r \dot{\phi} \underline{e}_\phi) - (2)$$

The relativistic momentum is:

$$\underline{p} = \gamma m \underline{\dot{r}} - (3)$$

where

$$\gamma = \left( m(r) - \frac{\underline{\dot{r}} \cdot \underline{\dot{r}}}{c^2} \right)^{-1/2} - (4)$$

By definition:

$$\underline{p} = \frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} - (5)$$

where  $\mathcal{L}$  is the Lagrangian. Choose:

$$\mathcal{L} = -mc^2 \left( m(r) - \frac{\underline{\dot{r}} \cdot \underline{\dot{r}}}{c^2} \right)^{1/2} + \frac{nmG}{r} - (6)$$

and use

$$U = -\frac{nmG}{r} - (7)$$

The Euler Lagrange equations are:

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} = \frac{\partial \mathcal{L}}{\partial \underline{r}} - (8)$$

and from Eq. (6):

$$\underline{p} = \frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} = \gamma m \underline{\dot{r}} - (9)$$

Q.E.D. So

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} \right) = m \frac{d\gamma}{dt} \underline{\dot{r}} + m\gamma \underline{\ddot{r}} - (10)$$

Note that:

$$\begin{aligned}\frac{\partial}{\partial \underline{r}} \left( \frac{1}{r} \right) &= \frac{\partial}{\partial \underline{r}} (\underline{r} \cdot \underline{r})^{-1/2} \\ &= 2 \underline{r} \left( -\frac{1}{2} (\underline{r} \cdot \underline{r})^{-3/2} \right) \quad - (11) \\ &= -\frac{\underline{r}}{(\underline{r} \cdot \underline{r})^{3/2}} = -\frac{\underline{r}}{r^3} = -\frac{1}{r^2} \underline{e}_r\end{aligned}$$

So:

$$m \frac{d\underline{v}}{dt} \underline{\dot{r}} + m \gamma \underline{\ddot{r}} = -\frac{m G}{r^2} \underline{e}_r \quad - (12)$$

like:

$$\underline{\dot{r}} = \frac{1}{m(r)^{1/2}} \left( \dot{r} \underline{e}_r + r \dot{\phi} \underline{e}_\phi \right) \quad - (13)$$

and

$$\underline{\ddot{r}} = \frac{1}{m(r)^{1/2}} \frac{d}{dt} \left( \dot{r} \underline{e}_r + r \dot{\phi} \underline{e}_\phi \right) \quad - (14)$$

here:

$$\underline{\dot{r}} = \frac{1}{m(r)^{1/2}} \left( \dot{r} \underline{e}_r + r \dot{\phi} \underline{e}_\phi \right) \quad - (15)$$

and

$$\underline{\ddot{r}} = \frac{1}{m(r)^{1/2}} \left( (\ddot{r} - r \dot{\phi}^2) \underline{e}_r + (r \ddot{\phi} + 2 \dot{r} \dot{\phi}) \underline{e}_\phi \right) \quad - (16)$$

It follows that:

$$\frac{1}{m(r)^{1/2}} \left( \frac{d\gamma}{dt} \underline{\dot{r}} + \gamma (\ddot{r} - r \dot{\phi}^2) \right) \quad - (17)$$

Let  $\Omega_r$  be seen as  $\frac{1}{r}$ , and

$$= -\frac{m G}{r} \left( \frac{1}{r} + \Omega_r \right)$$

$$\frac{1}{m(r)^{1/2}} \left( \frac{d\gamma}{dt} r \dot{\phi} + r \ddot{\phi} + 2 \dot{r} \dot{\phi} \right) = 0 \quad - (18)$$

By definition:

$$L = |\underline{r} \times \underline{p}| = \frac{\gamma m r^2 \dot{\phi}}{m(r)} \quad (11)$$

It follows from eq. (18) that:

$$\frac{dL}{dt} = 0 \quad (12)$$

The angular momentum is therefore defined as:

$$\underline{L} := \frac{1}{n(r)^{1/2}} \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} = \underline{r} \times \underline{p} \quad (13)$$

and the force equation is:

$$\underline{F} = \frac{d}{dt} (\gamma m \dot{\underline{r}}) = -\frac{n m G}{r^2} \underline{e}_r + \underline{\Omega} \times \underline{r} \quad (14)$$

In order to clarify the meaning of eq. (8), note that

$$\underline{p} = \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} = \frac{\gamma m}{n(r)^{1/2}} (\dot{\underline{r}} \underline{e}_r + r \dot{\phi} \underline{e}_\phi) \quad (15)$$

we:

$$\begin{aligned} \mathcal{L} &= \frac{n(r)^{1/2}}{2mc^2} \left( m(r) - \frac{\dot{\underline{r}} \cdot \dot{\underline{r}}}{n(r)c^2} \right)^{1/2} + \frac{n M G}{r} \\ &= \frac{n(r)^{1/2}}{2mc^2} \left( n(r) - \frac{\dot{\underline{r}} + r^2 \dot{\phi}^2}{n(r)c^2} \right)^{1/2} + \frac{n M G}{r} \end{aligned} \quad (16)$$

Therefore:

$$\frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} = \frac{\gamma m}{n(r)} (\dot{\underline{r}} \underline{e}_r + r \dot{\phi} \underline{e}_\phi) \quad (17)$$

$$\boxed{\frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} = \frac{\partial \mathcal{L}}{\partial \dot{r}} \underline{e}_r + \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \underline{e}_\phi} \quad (18)$$

Similarly:

$$\boxed{\frac{\partial \mathcal{L}}{\partial \underline{r}} = \frac{\partial \mathcal{L}}{\partial r} \underline{e}_r + \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \underline{e}_\phi} \quad - (19)$$

$$\underline{\quad} = \underline{\nabla} \mathcal{L}$$

The Euler Lagrange equation is therefore:

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} \right) = \underline{\nabla} \mathcal{L} \quad - (20)$$

From eq. (16):  $\underline{\nabla} \mathcal{L} = -\frac{m\Omega^2}{r} \underline{e}_r \quad - (21)$

Therefore  $\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} \right) = \frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{r}} \underline{e}_r \right) + \frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \underline{e}_\phi \right)$

$$= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}} \underline{e}_r + \frac{\partial \mathcal{L}}{\partial \dot{r}} \frac{d \underline{e}_r}{dt} + \frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \underline{e}_\phi + \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \frac{d \underline{e}_\phi}{dt} \quad - (22)$$

where:  $\frac{d \underline{e}_r}{dt} = \dot{\phi} \underline{e}_\phi, \quad \frac{d \underline{e}_\phi}{dt} = -\dot{\phi} \underline{e}_r \quad - (23)$

so  $\boxed{\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\underline{r}}} = \left( \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}} - \frac{\dot{\phi}}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \underline{e}_r + \left( \frac{\partial \mathcal{L}}{\partial \dot{r}} \dot{\phi} + \frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \right) \underline{e}_\phi} \quad - (24)$

So using the Lagrangian:

$$= -m(r)^{1/2} m c^2 \left( m(r) - \frac{\dot{r}^2 + r^2 \dot{\phi}^2}{m(r) c^2} \right)^{1/2} + \frac{n m \Gamma}{r} \quad (25)$$

we found that:

$$\frac{\partial \mathcal{L}}{\partial \dot{r}} = \frac{\gamma}{m(r)^{1/2}} \dot{r}, \quad \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \frac{\gamma}{m(r)^{1/2}} r^2 \dot{\phi} \quad (26)$$

and:

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}} = \frac{m}{m(r)^{1/2}} \left[ \frac{d}{dt} (\gamma \dot{r}) - \gamma r \dot{\phi}^2 \right] \underline{e}_r - (27)$$

$$+ \left[ \frac{\partial \mathcal{L}}{\partial \dot{r}} \dot{\phi} + \frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \right] \underline{e}_\phi$$

$$= \frac{m}{m(r)^{1/2}} \left( \frac{d\gamma}{dt} \dot{r} \underline{e}_r + \gamma (\ddot{r} - r \dot{\phi}^2) \underline{e}_r \right. \\ \left. + \gamma \dot{r} \dot{\phi} \underline{e}_\phi + \frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) \underline{e}_\phi \right)$$

Finally we:

$$\frac{d}{dt} \left( \frac{1}{r} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) = \frac{1}{m(r)^{1/2}} \frac{d}{dt} (\gamma r \dot{\phi}) \quad (28)$$

$$= \frac{1}{m(r)^{1/2}} \left( \frac{d\gamma}{dt} r \dot{\phi} + \gamma (\dot{r} \dot{\phi} + r \ddot{\phi}) \right)$$

we found that:

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}} = \frac{1}{m(r)^{1/2}} \left[ \frac{d\gamma}{dt} (\dot{r} \underline{e}_r + r \dot{\phi} \underline{e}_\phi) \right. \quad (29)$$

$$\left. + \gamma ((\ddot{r} - r \dot{\phi}^2) \underline{e}_r + (r \ddot{\phi} + 2 \dot{r} \dot{\phi}) \underline{e}_\phi) \right]$$

) which is eq. (28) of N. To 415(4) Q.E.D.

With the Lagrangian (25) the fermionic and  
gravitino methods give exactly the same results.  
The double cross checked equations of motion are

Therefore:

$$\frac{dV}{dt} \dot{r} + V(\ddot{r} - r\dot{\phi}^2) = -m(r) \frac{MG}{r} \left( \frac{1}{r} + 2r \right) - (30)$$

and

$$\frac{dV}{dt} r\dot{\phi} + r\ddot{\phi} + 2\dot{r}\dot{\phi} = 0 - (31)$$

The conserved angular momentum is:

$$L := \frac{\gamma m r^2}{m(r)} \dot{\phi} - (32)$$

and Eq. (31) imply:

$$\frac{dL}{dt} = 0 - (33)$$

---